

On Parametrized KAM Theory

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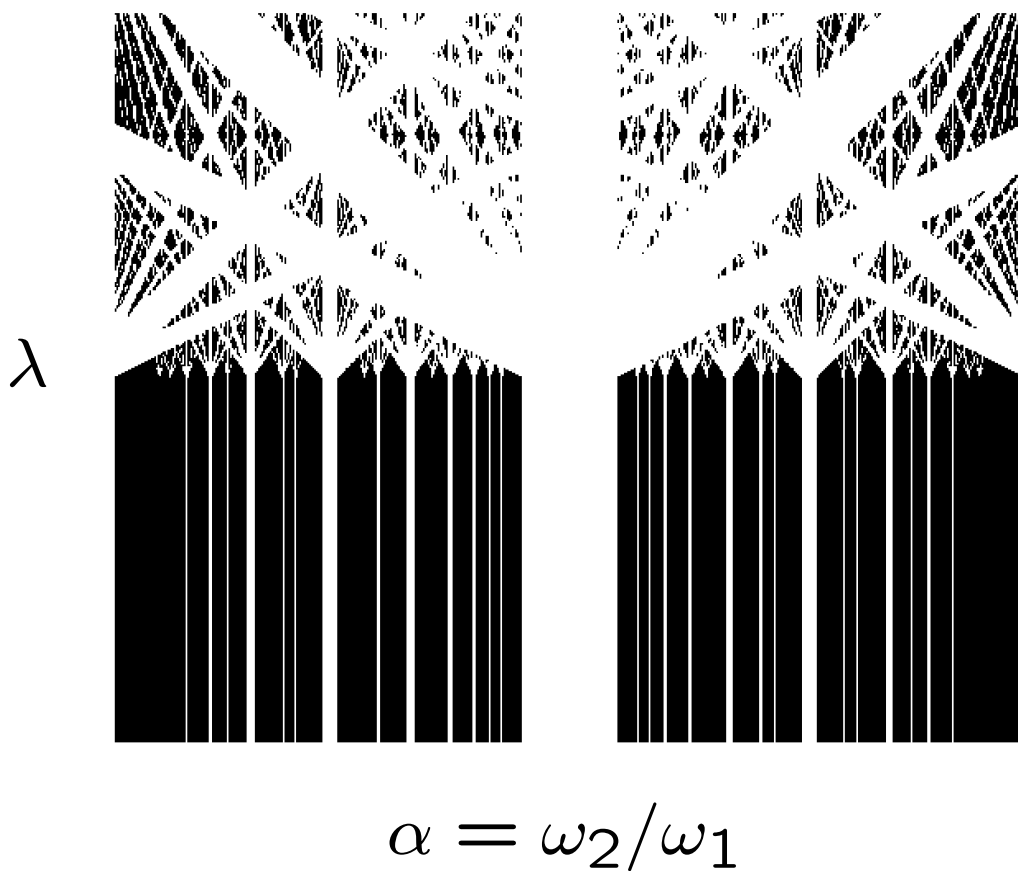
AIMS Dresden, May 2010

Motivation

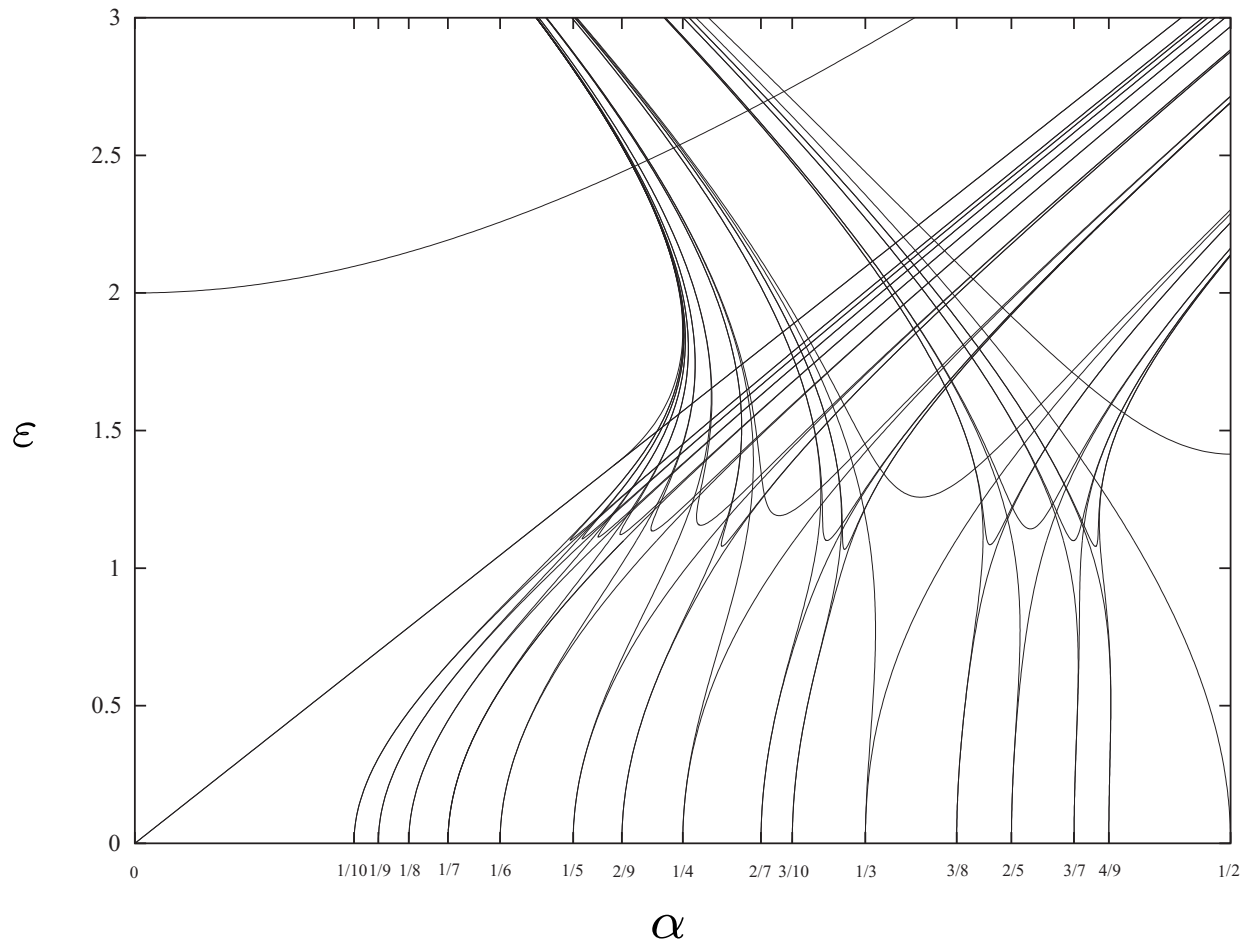
Parametrized KAM Theory and Quasi-Periodic Bifurcation Theory developed in:

Moser (1966, 1967), Chenciner-Iooss (1979), Chenciner (1985), Braaksma-B (1987), Huitema (1988), B-Huitema-Takens-Braaksma (1990), B-Huitema (1995), B-Huitema-Sevryuk (1996), Hanßmann (1998, 2006, 2006), B-Hanssmann-Jorba-Villanueva-Wagener (2003), Hoo (2005), Wagener (2002 and to appear), B-Hoo-Naudot (2007), B-Hanßmann-You (2005, 2006 and in preparation), B-Hanßmann-Hoo (2007), B-Ciocci-Hanßmann (2007), B-Ciocci-Hanßmann-Vanderbauwhede (2009), B-Takens (2009), B-Sevryuk (to appear), B-Hanßmann-Wagener (in preparation)

Marriage of KAM– and Singularity Theory \rightsquigarrow
Cantorisation of semi-algebraic stratifications
(or Cantor stratifications)



Cantorized fold in quasi-periodic bifo's. E.g. in q-p (Hamiltonian) center-saddle bifo and in 'Ruelle-Takens' bifo of (dissipative) maps from invariant circles to invariant 2-tori. Gaps (white) contain 'interesting' dynamics, like periodicity and chaos, generically coexisting with positive measure of quasi-periodicity.



Resonance tongues in the Arnol'd family of circle maps

$$A_{\alpha,\varepsilon}(x) = x + 2\pi\alpha + \varepsilon \sin x$$

for $\varepsilon < 1$ circle diffeomorphisms associated to the lower half of the previous figure

Parametrized KAM Theory (dissipative case)

Phase space $M = \mathbb{T}^n \times \mathbb{R}^m = \{(x, z)\}$
parameter space $P \subseteq \mathbb{R}^s = \{\mu\}$ (open)

Starting point:

integrable C^∞ -family $X = X_\mu(x, z)$

$$\begin{aligned}\dot{x} &= \omega(\mu) + O(|z|) \\ \dot{z} &= \Omega(\mu)z + O(|z|^2),\end{aligned}$$

with O -estimates (locally) uniform in μ ,
 $\omega(\mu) \in \mathbb{R}^n$ and $\Omega(\mu) \in gl(m, \mathbb{R})$

Integrability of $X \Leftrightarrow x$ -independence

Interest: persistence properties of the family
 $T_\mu = \mathbb{T}^n \times \{0\}$ of invariant n -tori with parallel
dynamics

BHT nondegeneracy (1990)

Nondegeneracy condition on product map

$$\omega \times \Omega : P \rightarrow \mathbb{R}^n \times gl(m, \mathbb{R})$$

at $\mu = \mu_0 \in P$: ‘simultaneously’

1. $\mu \in \mathbb{R}^s \mapsto \omega(\mu) \in \mathbb{R}^n$ submersion
(think of Kolmogorov-nondegeneracy)

2. $\mu \in \mathbb{R}^s \mapsto \Omega(\mu) \in gl(m, \mathbb{R})$ versal
unfolding of $\Omega(\mu_0)$

i.e., transversal to orbit of $\Omega(\mu_0)$
under the adjoint action of $GL(m, \mathbb{R})$

Assumption throughout: $\Omega(\mu_0)$ invertible
(relaxed by Wagener (to appear)
and B-Hanßmann-Wagener (in preparation))

Case: $\Omega(\mu_0)$ simple eigenvalues

Eigenvalues $\Omega(\mu)$ (where $N_1 + 2N_2 = m$):

$$(\delta_1, \dots, \delta_{N_1}, \alpha_1 \pm i\beta_1, \dots, \alpha_{N_2} \pm i\beta_{N_2})$$

with $\beta_j > 0$ for $1 \leq j \leq N_2$

Call $\beta = (\beta_1, \beta_2, \dots, \beta_{N_2})$ *normal frequencies*

By simpleness of eigenvalues, map

$$\text{spec} : gl(m, \mathbb{R}) \rightarrow \mathbb{R}^{N_1} \times \mathbb{R}^{N_2} \times \mathbb{R}^{N_2}; \quad \Omega \mapsto (\delta, \alpha, \beta),$$

parametrizes $GL(m, \mathbb{R})$ orbit space near $\Omega(\mu_0)$

BHT nondegeneracy condition means that

$$P \ni \mu \mapsto (\omega \times (\text{spec} \circ \Omega))(\mu) \in \mathbb{R}^n \times \mathbb{R}^{N_1} \times \mathbb{R}^{N_2} \times \mathbb{R}^{N_2}$$

is submersion (at $\mu = \mu_0$) !!

Inverse Function Theorem

\leadsto 'reparametrization'

$$\mu \leftrightarrow (\omega, \delta, \alpha, \beta)$$

on $A \subseteq P$ (open)

Diophantine conditions

For $\tau > n-1$ and $\gamma > 0$ define (τ, γ) -Diophantine normal-internal frequency vectors by

$$D_{\tau, \gamma}(\mathbb{R}^n; \mathbb{R}^{N_2}) = \{(\omega, \beta) \in \mathbb{R}^n \times \mathbb{R}^{N_2} \mid |\langle \omega, k \rangle + \langle \beta, \ell \rangle| \geq \gamma |k|^{-\tau}, \\ \forall k \in \mathbb{Z}^n \setminus \{0\} \text{ and } \forall \ell \in \mathbb{Z}^{N_2} \text{ with } |\ell| \leq 2\}$$

Nowhere dense, positive measure, closed half line property

Defining $\Gamma = (\omega \times (\text{spec} \circ \Omega))(A)$, w.l.o.g. assume that Γ has 'product form'

$$\Gamma = \Gamma_\omega \times \Gamma_\delta \times \Gamma_\alpha \times \Gamma_\beta$$

Shrunken version

$$\Gamma_\gamma = \{(\omega, \delta, \alpha, \beta) \in \Gamma \mid \text{dist}((\omega, \delta, \alpha, \beta), \partial\Gamma) \geq \gamma\}$$

of Γ as well as

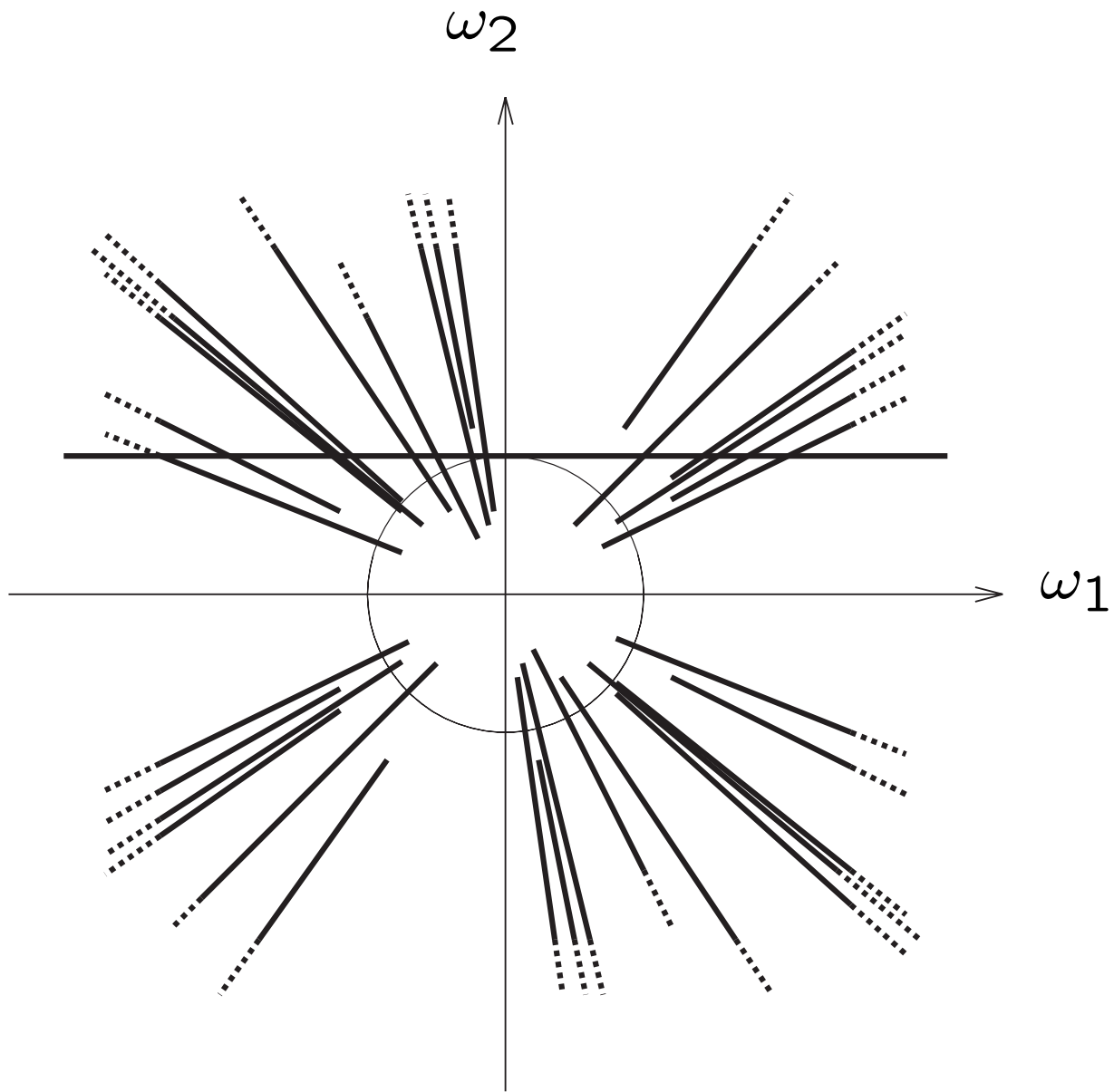
$$D_{\tau, \gamma}(\Gamma_\gamma) = \Gamma_\gamma \cap (D_{\tau, \gamma}(\mathbb{R}^n; \mathbb{R}^{N_2}) \times \Gamma_\delta \times \Gamma_\alpha)$$

and $D_{\tau, \gamma}(A_\gamma) \subset A$

Closed half lines turn into closed linear half spaces of dimension $1 + N_1 + N_2$

KAM theory: this geometry, up to diffeo's, is inherited by perturbations

!!



Diophantine set $D_{\tau, \gamma}(\mathbb{R}^n)$ has closed half line geometry For $\tau > n - 1$ intersection $D_{\tau, \gamma}(\mathbb{R}^n) \cap \mathbb{S}^{n-1}$ is Cantor set; measure complement $\mathbb{S}^{n-1} \setminus D_{\tau, \gamma}(\mathbb{R}^n) = O(\gamma)$ as $\gamma \downarrow 0$

Parametrized KAM ('dissipative')

Thm. Let integrable C^∞ -family $X = X_\mu(x, z)$ be BHT nondegenerate on $\mathbb{T}^n \times \{0\} \times A$

Assume $\det \Omega(\mu) \neq 0$

Then, for $0 < \gamma \ll 1$, $\exists C^\infty$ -nbhd of X , such that, $\forall \tilde{X} \in \mathcal{O}$, \exists (local) C^∞ -map $\Phi : \mathbb{T}^n \times \mathbb{R}^m \times A \rightarrow \mathbb{T}^n \times \mathbb{R}^m \times A$, with:

1. Φ is a C^∞ -near identity diffeo preserving projections to the parameter space P
2. Image $\Phi(\mathbb{T}^n \times \{0\} \times D_{\tau, \gamma}(A_\gamma))$ is \tilde{X} -invariant, and restriction $\widehat{\Phi} = \Phi|_{\mathbb{T}^n \times \{0\} \times D_{\tau, \gamma}(A_\gamma)}$ conjugates X to \tilde{X} , that is

$$\widehat{\Phi}_* X = \tilde{X}$$

3. Φ preserves normal linear behavior of the tori T_μ for $\mu \in D_{\tau, \gamma}(A_\gamma)$

Discussion I

1. Family X is quasi-periodically stable on torus union $\mathbb{T}^n \times \{0\} \times D_{\tau, \gamma}(A_\gamma)$
(Item 3 \sim *normal linear stability*)

2. Preservation of structures allowed:
Hamiltonian, volume preserving,
equivariant and reversible,
with or without external parameters,
combinations of the above
Axiomatic: *admissible* structures

3. Simpleness of eigenvalues of $\Omega(\mu)$ can be dropped in full 'admissible' generality

Arnold's theory of matrices depending on parameters carries over: Normalization to *Linear Centralizer Unfoldings*

Example in normal 1 : -1 -resonance

Normally trivial cases

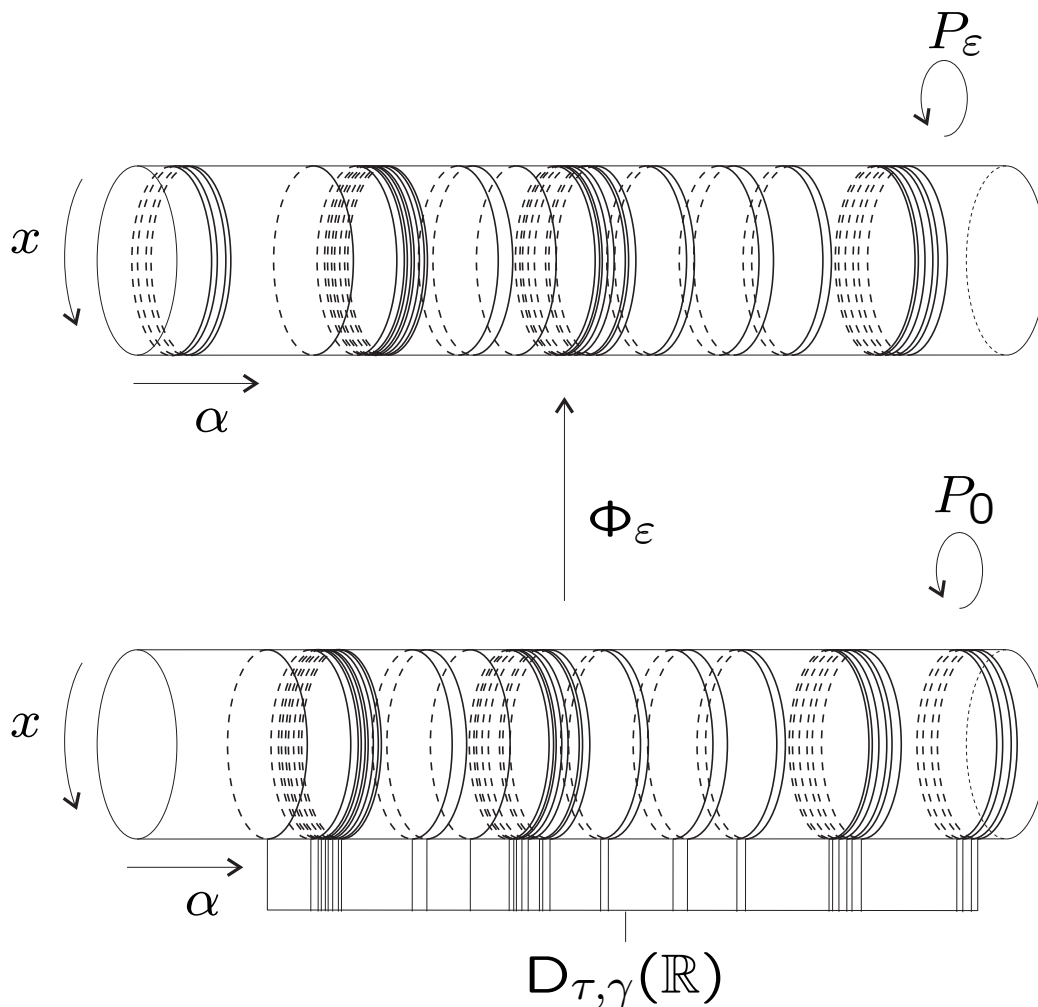
In these cases $m = 0$ ($\Rightarrow \Omega(\mu) \equiv 0$):

1. Normally Hyperbolic Dissipative KAM
(families of quasi-periodic attractors)
2. 'Classical' Lagrangean KAM (Hamiltonian,
a similar reversible case exists too)
 $X(x, y) = \omega(y) \frac{\partial}{\partial x}$,
Compare with proof Pöschel (1982)
3. KAM for codimension 1 tori in volume
preserving setting
4. &c.

Latter two *direct* consequence by
'localisation': invariant torus given by $y = \nu$,
then define $y_{loc} = y - \nu$ and pass to

$$X_{loc}(x, y_{loc}, \nu) := X(x, y_{loc} + \nu)$$

NB: ν is *distinguished* parameter



Family of quasi-periodic attractors and quasi-periodic stability: simplest case for circle maps

$$P_{\epsilon; \alpha}(x) = x + 2\pi\alpha + \epsilon f(x, \epsilon; \alpha)$$

Isotropic Hamiltonian tori

Slightly adapted format family X valid for Hamiltonian subtori:

- $x \in \mathbb{T}^n, y \in \mathbb{R}^n, z \in \mathbb{R}^m$ with m even and symplectic form $dx \wedge dy + dz^2$

$$X(x, y, z) = \omega(y) \frac{\partial}{\partial x} + \Omega(y) z \frac{\partial}{\partial z}$$

- 'Actions' y distinguished parameters
Compare with proof Pöschel (1989)
- Analogous for reversible setting

Idea also used by Herman (lecture 1990's)
compare with [BHS (1996), BS (to appear)]

On the Axioms

Given: Pair (G, V) with

$G \neq \{I_m\}$ Lie subgroup of $GL(m, \mathbb{R})$

$V \neq \{0\}$ linear subspace of $gl(m, \mathbb{R})$ such that

$$Q \in G, \&\Omega \in V \Rightarrow Q\Omega Q^{-1} \in V$$

(invariance under G -action)

G -orbit: $O(\Omega, G) = \{Q\Omega Q^{-1} : Q \in G\} \subseteq V$

$O(\Omega, G)$ is smooth manifold, e.g., if

if G is semi-algebraic subset (**Axiom A1**)

Also assume: $\Omega \in V \Rightarrow \Omega^T \in V$ (**Axiom A2**)

Pair (G, V) called *linear structure*

NB: **Axiom A3** later

Applies to $G = GL(m, \mathbb{R}), SL(m, \mathbb{R}), SP(m, \mathbb{R})$

or $GL_{+R}(m, \mathbb{R})$ with V their Lie algebra,

and to reversible case

$(G, V) = (GL_{+R}(m, \mathbb{R}), gl_{-R}(m, \mathbb{R}))$,

with $R \in GL(m, \mathbb{R})$ an involution, where

$$gl_{-R}(m, \mathbb{R}) := \{A \in gl(m, \mathbb{R}) \mid AR = -RA\}$$

Unfoldings within (G, V)

Let $\Pi_p(V)$ be the set of all smooth p -parameter families on V (i.e., germs at $\mu = 0$)

Unfolding $\Omega \in \Pi_p(V)$ of Ω_0 *versal* if
 $\forall B \in \Pi_q(V)$ with $B(0) = \Omega_0$, \exists local map
 $(\rho, Q) : (\mathbb{R}^q, 0) \rightarrow (\mathbb{R}^p, 0) \times (G, \mathcal{I}_m)$ such that

$$B(\mu) = Q(\mu)\Omega(\rho(\mu))Q^{-1}(\mu)$$

Thm. [BHN] Family Ω is *versal* at $\mu = 0 \iff$
 $D_0\Omega(\mathbb{R}^p) + T_{\Omega_0}(O(\Omega_0, G)) = V$ (transversal)

Minimal p called *codimension*

\rightsquigarrow *universal* unfolding

NB: Invariance of eigenvalues:

\rightsquigarrow always $\text{cod } \Omega_0 > 0$

Lemma. $T_{\Omega_0}(O(\Omega_0, G)) = \text{ad}\Omega_0(\mathfrak{g})$, with \mathfrak{g}
Lie algebra of G and $\text{ad}\Omega_0(Y) = [\Omega_0, Y]$

Versal Unfolding Stability Thm

Linear unfolding $\Omega \in \Pi_p(V)$ of Ω_0 is called Linear Centralizer Unfolding (LCU) if $p = \text{cod}\Omega_0$ and if

$$D_0\Omega(\mathbb{R}^p) = \ker \text{ad } \Omega_0^T \subseteq V$$

NB: $\ker \text{ad } \Omega_0^T \oplus \text{im ad } \Omega_0 = V$

Thm. $\Omega \in \Pi_p(V)$ versal within linear structure (G, V) . Then $\exists C^1$ -neighbourhood \mathcal{V} of Ω in $\Pi_p(V)$, and $\forall B \in \mathcal{V}$, \exists map (ρ, Q) , such that

i. $B(\mu) = Q(\mu)\Omega(\rho(\mu))Q^{-1}(\mu)$, for small μ ;

ii. Reparametrization ρ is C^1 -near identity

Proof: IFT

Linear stability of equilibria (dissipative case)

Given: family $X_\mu(z) = [\Omega(\mu)z + O(|z|^2)]\frac{\partial}{\partial z}$
 $\Omega(\mu) \in gl(m, \mathbb{R}), \mu \in P$

Assume $\Omega_0 = \Omega(0)$ invertible and $\mu \mapsto \Omega(\mu)$
versal (nondegeneracy)

Zero set $\mathcal{Z}(X) \subset \mathbb{R}^m \times P$ of family X :

$\mathcal{Z}(X) = \{z = 0, \mu \text{ in full nbhd of } 0 \in P\}$

Lemma. $\forall C^2$ -small perturbation \tilde{X} of X , \exists
diffeomorphism $\Phi : \mathbb{R}^m \times P \rightarrow \mathbb{R}^m \times P$
(of form $\Phi(z, \mu) = (\phi_\mu(z), \rho(\mu))$), with

i. Φ is C^1 -near identity;

ii. $\Phi(\mathcal{Z}(X)) = \mathcal{Z}(\tilde{X})$;

iii. Φ preserves the linear behaviour of X

Discussion II

Conclusion of Lemma \longleftrightarrow *linear stability* of X at equilibrium $(z, \mu) = (0, 0)$

By IFT $\mathcal{Z}(\tilde{X})$ locally is smooth graph $(z(\mu), \mu)$ in $\mathbb{R}^m \times P$ over P

Trivial application to Parametrized KAM Theory when perturbation \tilde{X} of integrable family X *also* is integrable !

Full KAM Thm as before. Linear stability solves small divisor problems. Newtonian iteration with *homological equation*

Similar statements for other linear structures: e.g., symplectic, volume preserving, equivariant, reversible, etc.

Proof: Moser-Pöschel-Huitema-Hoo [BHN (2007), BCHV (2009)]

Admissibility: linear structure Axioms A1, A2 PLUS **Axiom A3**: ‘homogeneity’ (Taylor-Fourier truncations preserve structure)

Normal 1 : -1 -resonance I

B, Hoo, Ciocci, Hanßmann, Naudot,
Vanderbauwhede (2007-09)

$\mathbb{R}^4 = \{z_1, z_2, z_3, z_4\}$ with symplectic form
 $\sigma = dz_1 \wedge dz_3 + dz_2 \wedge dz_4$

Let Y be Hamiltonian vector field given by

$$Y(z) = \Omega_0 z \frac{\partial}{\partial z} \in \mathcal{X}^\sigma(\mathbb{R}^4),$$

where $\Omega_0 \in sp(4, \mathbb{R})$

In $1 : -1$ resonance if symplectically similar
to

$$N_\varepsilon = \begin{pmatrix} 0 & -\lambda_0 & 0 & 0 \\ \lambda_0 & 0 & 0 & 0 \\ \varepsilon & 0 & 0 & -\lambda_0 \\ 0 & \varepsilon & \lambda_0 & 0 \end{pmatrix},$$

where $\lambda_0 \neq 0$ and $\varepsilon = 0, \pm 1$

Generic case: $\varepsilon = \pm 1$, just take $\varepsilon = 1$

Normal 1 : -1 -resonance II

For $\varepsilon = 1$ (generic case),
an LCU of N_ε is given by family

$$\Omega(\mu) = \begin{pmatrix} 0 & -\lambda_0 - \mu_1 & -\mu_2 & 0 \\ \lambda_0 + \mu_1 & 0 & 0 & -\mu_2 \\ 1 & 0 & 0 & -\lambda_0 - \mu_1 \\ 0 & 1 & \lambda_0 + \mu_1 & 0 \end{pmatrix}$$

where $\mu_1, \mu_2 \in \mathbb{R}$ (codimension = 2)

Eigenvalues $\Omega(\mu)$:

$$\pm(i(\lambda_0 + \mu_1) \pm \sqrt{\mu_2}),$$

1. Elliptic for $\mu_2 < 0$
2. Parabolic (or Doubly Elliptic) for $\mu_2 = 0$:
 $\rightsquigarrow i(\lambda_0 - \mu_1), -i(\lambda_0 - \mu_1)$
3. Hyperbolic for $\mu_2 > 0$

Similar for the semisimple case: cod = 4

Similar for the reversible case:

same codimensions

Normal 1 : -1 -resonance III

Recall: eigenvalues of $\Omega(\mu)$

$$\pm i(\lambda_0 + \mu_1 \pm \sqrt{\mu_2})$$

Normal frequencies of X given by

$$\mu_2 \leq 0 \Rightarrow$$

$$\omega^N(\mu) = (\lambda_0 + \mu_1, \lambda_0 + \mu_1);$$

$$\mu_2 > 0 \Rightarrow$$

$$\omega^N(\mu) = (\lambda_0 + \mu_1 + \sqrt{\mu_2}, \lambda_0 + \mu_1 - \sqrt{\mu_2})$$

(Generalized) frequency map \mathcal{F} given by:

$$\mathcal{F}(\omega, \mu) = \begin{cases} (\omega, \lambda_0 + \mu_1, \lambda_0 + \mu_1) & \text{for } \mu_2 \leq 0 \\ (\omega, \lambda_0 + \mu_1 + \sqrt{\mu_2}, \lambda_0 + \mu_1 - \sqrt{\mu_2}) & \\ & \text{for } \mu_2 > 0. \end{cases}$$

Corresponding Diophantine conditions

$$|\langle \omega, k \rangle + \ell(\lambda_0 + \mu_1)| \geq \gamma |k|^{-\tau}$$

$$|\langle \omega, k \rangle + \ell_1(\lambda_0 + \mu_1) + \ell_2 \sqrt{|\mu_2|}| \geq \gamma |k|^{-\tau},$$

$\forall k \in \mathbb{Z}^n \setminus \{0\}$ and $\forall |\ell| = 0, 1, 2 \quad \forall |\ell_1|, |\ell_2| \leq 2$
with $(\ell_1 \pm \ell_2)$ even

Nonlinear aspects I

Quasi-periodic Hamiltonian Hopf bifurcation

Higher order terms (VdMeer):

$$\hat{H} = (\nu_1(\mu) + \lambda_0)S + N + \nu_2(\mu)M + \frac{1}{2}b(\mu)M^2 + c_1(\mu)SM + c_2(\mu)S^2 + \dots$$

$$M = \frac{1}{2}(z_1^2 + z_2^2), N = \frac{1}{2}(z_1^3 + z_2^4), S = z_1z_4 - z_2z_3$$

Rotational symmetry

Integrable (swallowtail) versus

nearly-integrable (Cantorized swallowtail)

$$H(x, y, z, \mu) = \langle \omega(\mu), y \rangle + \frac{1}{2} \langle Jz, \Omega(\mu)z \rangle + \dots$$

Nearly integrable KAM theory:

\rightsquigarrow preservation of *most degenerate* m -torus, !!
with hyperbolic and elliptic m -tori along crease
and thread

elliptic $(m + 1)$ -tori on 2-dimensional sheet
& *Lagrangian* $(m + 2)$ -tori above sheet
(both by 'standard' KAM)

Nonlinear aspects II

Quasi-periodically driven Lagrange top
(various devices)

Cushman-VdMeer
B-Hanßmann-Hoo (2007)

NB: All parameters 'distinguished':
~> weak stability
(only control frequency ratios)

Similarly: qp reversible Hopf

Knobloch-Vanderbauwhede
B-Ciocci-Hanßmann (2007)

Fewer parameters

More or less direct consequences of Parametrized KAM Theory:

1. Submersivity of frequency ratio

$$[\omega_1 : \omega_2 : \dots : \omega_n] \in \mathbb{P}^{n-1}(\mathbb{R})$$

\rightsquigarrow *weak* quasi-periodic stability
(all admissible structures)

2. Related to isoenergetic KAM
in Hamiltonian systems

B-Huitema (1991)

3. Rüssmann nondegeneracy too:

M.R. Herman lectures 1990's

B-Huitema-Sevryuk (1996), B-Sevryuk (to appear)

Discussion III

GAMES:

'Compensate' frequencies +
unfolding parameters
by 'distinguished' ones or by 'coefficients'

Play effectively with $\gamma = \gamma(\mu)$:
as small 'as perturbation allows'

CONCLUSIONS:

Cantor stratifications and Measure Theory:
many Hausdorff density points in appropriate
dimension (e.g., exponential condensation)
see [B-Sevryuk (1996, to appear)]
for references

EXAMPLES (in various admissible structures):

Quasi-periodic BIFO's (Hopf, period doubling)
With $\det \Omega_0 = 0$: qp saddle-node, cuspoids,
saddle-center, normally parabolic

(e.g., $H_\mu(x, y, t) = \frac{1}{2}y^2 + V_\mu(x) + \varepsilon f(x, \mu, t)$)
and umbilic cases [B-Hanßmann-You (2005)
and (2006)] Destruction of Lagrangean tori
by internal resonance [BHY (to appear)]

Global problems B-Cushman Fassò-Takens (2007)